

Impact of Baryon anti-Baryon annihilation on apparent strangeness enhancement in Λ/\bar{p} at SPS energy

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ICPAQGP 2023, Puri, Odisha

Outline

- ✓ Strangeness enhancement
- ✓ Measures of strangeness enhancement.
- ✓ Importance of anti-lambda to anti-proton ratio (Λ/\bar{p})
- ✓ Effects of hadronic interactions importantly baryon-anti- baryon annihilation on Λ , Λ/\bar{p} .

Baryon-anti- baryon ($B\bar{B}$) annihilation effect on

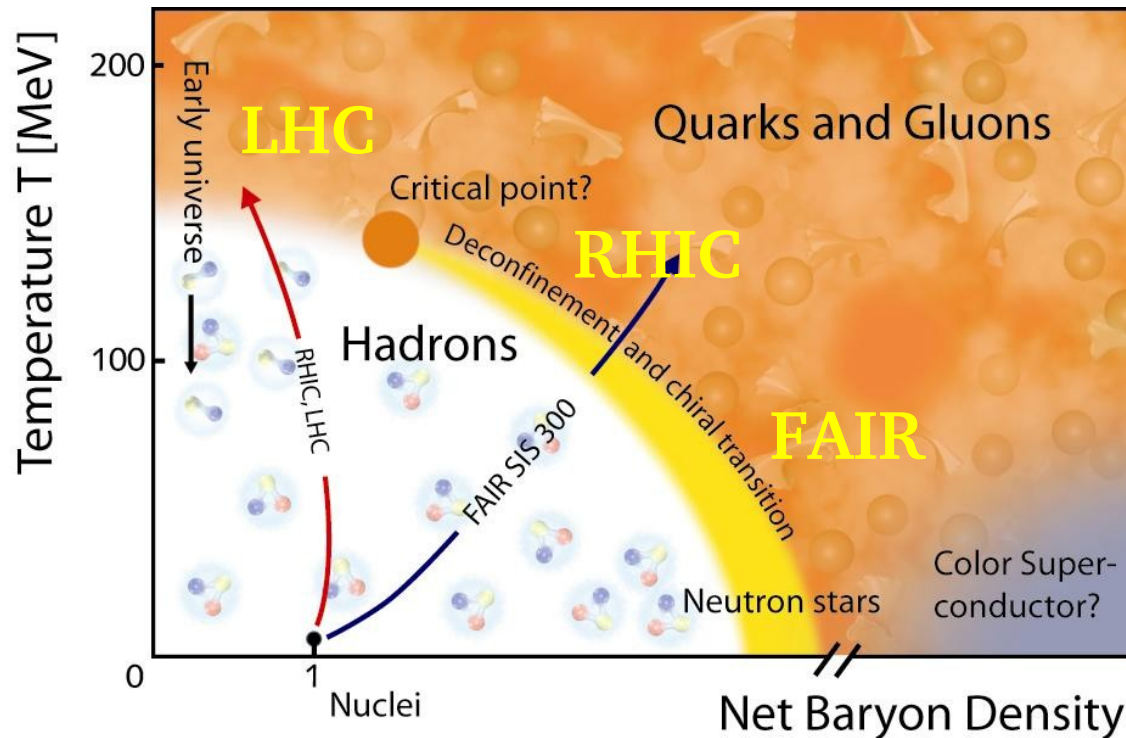
- $\langle m_T \rangle$ of Λ , Λ
- $\langle m_T \rangle$ with Centrality of Λ , Λ
- Yield Λ , \bar{p} w.r.to p_T and rapidity
- Annihilation fraction with beam energy
- Λ/\bar{p} Ratio

- ✓ Summary

Ref: Ekata Nandy and Subhasis Chattopadhyay,
European Physical Journal A, 58 (2022) 10, 199

Phase diagram of nuclear matter

- Under extreme conditions of temperature or pressure normal nuclear matter (hadronic phase) is likely to undergo a deconfinement phase transition to a quark gluon plasma phase.
- Effort to locate the phase boundary at different regions using different experiments.



LHC – Large Hadron Collider
($E_{\text{cm}} = 2.76\text{TeV} - 5.02\text{ TeV}$)

RHIC BES – Relativistic Heavy Ion Collider Beam Energy Scan
($E_{\text{cm}} = 7.7\text{ GeV} - 200\text{ GeV}$)

FAIR – Facility for Antiproton & Ion Research
($E_{\text{cm}} = 3\text{ GeV} - 9\text{ GeV}$)

Signatures of QGP

There is no unique signal that will identify QGP. Different signatures are used to search for QGP.

- J/ψ suppression
- Strangeness enhancement**
- Jet quenching
- Dilepton production

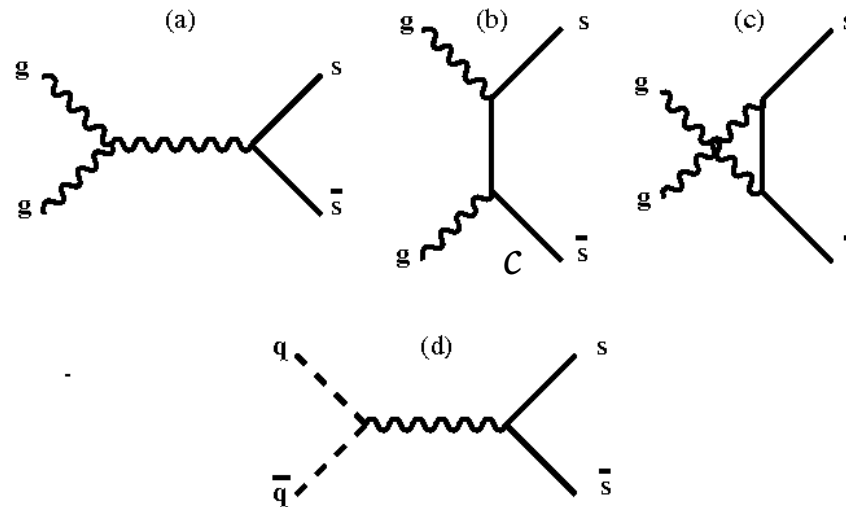
Strangeness production

*J. Rafelski and B. Müller first predicted Strangeness enhancement as a signature of deconfinement
As there is no initial valence strange quark, it produces from the reactions only.*

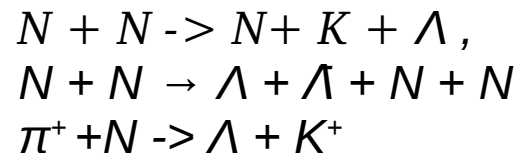
*Partonic
channel*

$$g + g \rightarrow ss$$

$$q + q \rightarrow ss$$



*Hadronic
channel*



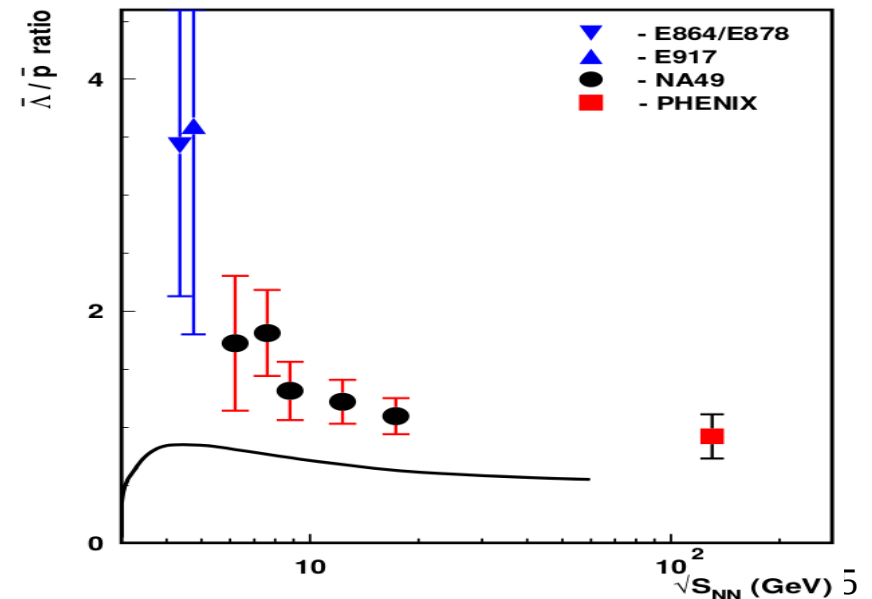
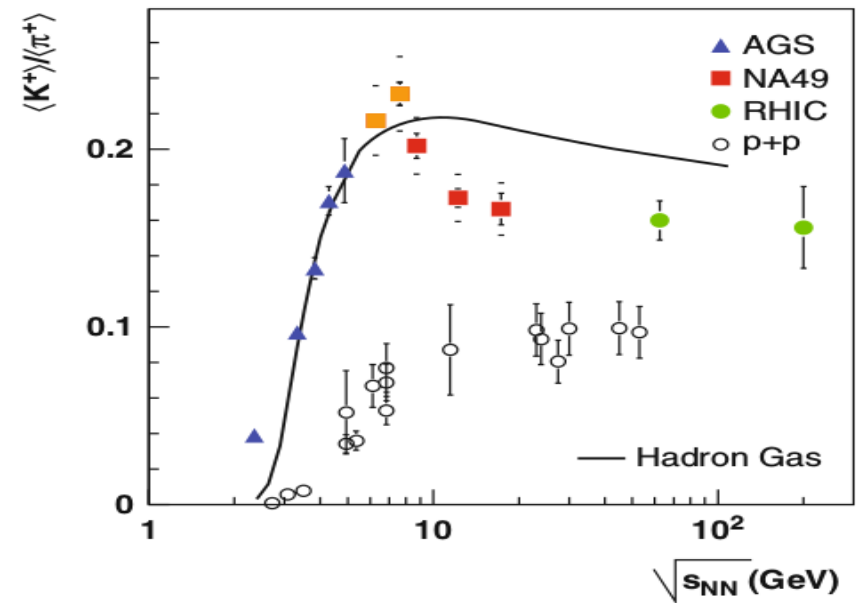
Why do we expect strangeness enhancement at low energy?

(Fermi Energy and Pauli Blocking)

- Because of higher abundance light quarks (u,d) in the medium they fill up the available low energy levels upto the fermi energy. Thus to produce a uu pair , required energy = fermi energy + $2m_u$
- Thus it is energetically favourable to produce ss pairs that require a threshold energy just double the mass of strange quark only.

Measures of strangeness enhancement

- ✚ A large enhancement in strange hadrons production relative to pp interaction was first reported at CERN-SPS.
- ✚ Enhanceent factor (relative to pp) $\left. \frac{dN(Pb + Pb)}{dy} \right|_{y=0} / \left. \frac{dN(p + p)}{dy} \right|_{y=0}$
- ✚ Interesting structures were observed in the strange-to-non-strange particle ratios.
- ✚ Non-monotonic variation of k/π as a function of collision energy was first observed at SPS.
- ✚ Similar behaviour was also observed in the baryon sector (Λ/\bar{p}), although with large uncertainty.
- ✚ Such non-monotonic variation is often attributed to the onset of partonic deconfinement.
- ✚ However hadronic interactions such as BB annihilation has a dominant role to play in strangeness baryon production & enhancement in the baryon sector .



Motivation of work

- Understand the contribution of hadronic interactions to the measures of strangeness enhancement.
- Final yields of baryons are highly sensitive to hadronic interactions at later stages of the collisions mainly from the baryon-anti baryon annihilation in a baryon rich environment (FAIR & SPS energies).
- So depending on the different annihilation cross-section of \bar{p} and Λ , this ratio (Λ/\bar{p}) may enhance.
- This study further aims to address whether the enhancement in the ratio (Λ/\bar{p}) can be explained from the consequence of hadronic interactions alone ?

Why Λ/\bar{p} ?

*Since anti-particles comprise of quarks produced in the reactions only,
they are regarded as a cleaner channel to probe strangeness
enhancement*

Details of model simulation

UrQMD (Ultra Relativistic Quantum Molecular Dynamics)

- Hadronic model, describe the phenomenology of particle production in pp, p+A and A+A collisions over a broad energy range
- The underlying degrees of freedom in UrQMD are hadrons and strings.
- Here an individual particle propagates on a straight line until the relative distance between two particles is smaller than the total interaction cross-section between two particles.
- The particle production dynamics is either governed by the decays of baryon or meson resonances or via a string excitation and fragmentation.

System : Pb+Pb/Au+Au

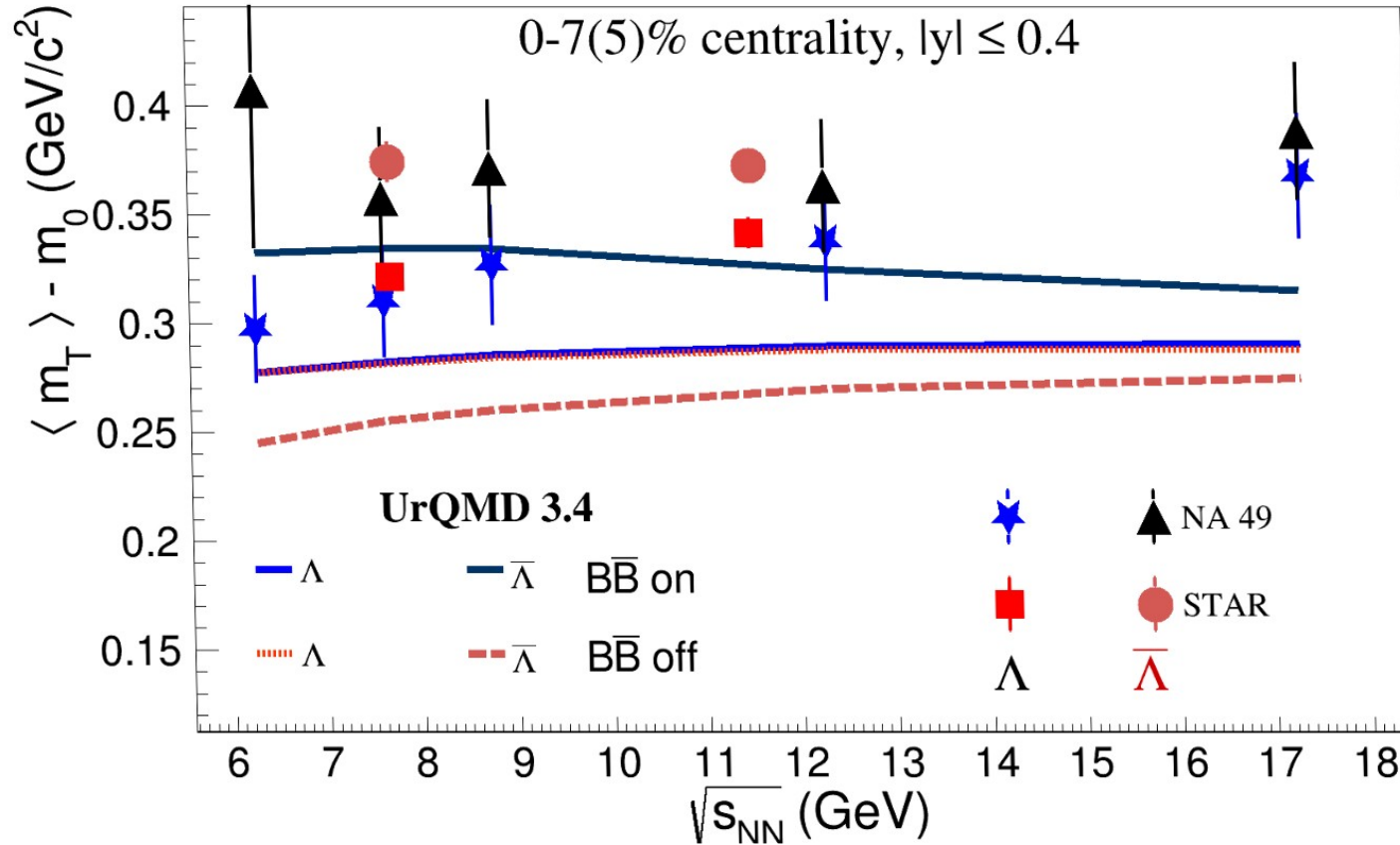
With & w/o incorporating baryon -anti baryon annihilation

Energy : 6.27 GeV, 7.62 GeV, 8.77 GeV, 12.3 GeV, 17.3 GeV

Centrality = 0-7%

Observables : Λ , \bar{p}

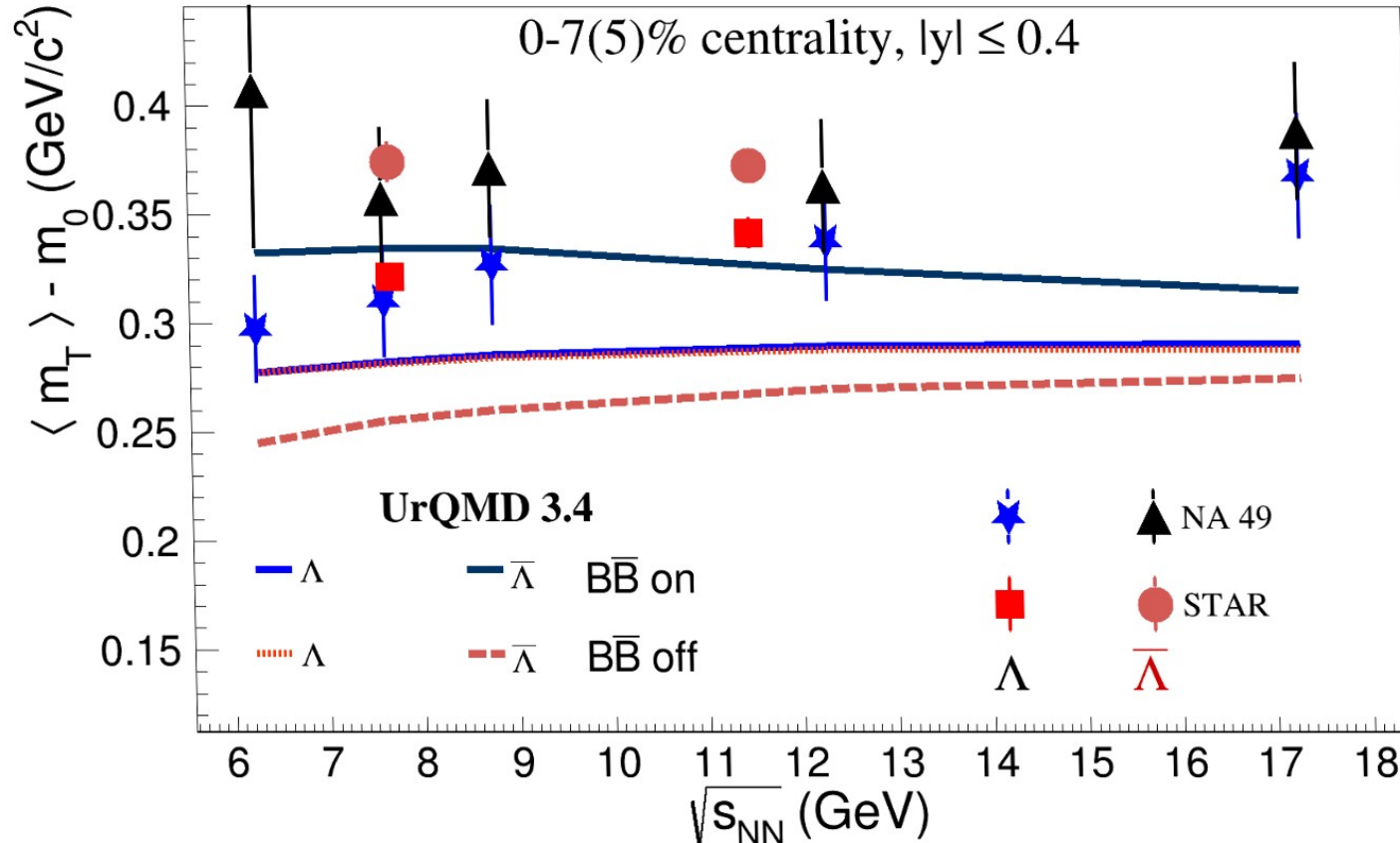
$B\bar{B}$ annihilation effect on $\Lambda, \bar{\Lambda} \langle m_T \rangle$



- ✓ Average transverse mass \equiv effective temperature. (T_{eff}) . $T_{\text{eff}} = T_{\text{th}} + 1/2 m\beta^2$
- ✓ Although Λ and $\bar{\Lambda}$ have same mass, their $\langle m_T \rangle$ are different in magnitude and the difference is seen to increase with the decrease in beam energy or increase in net baryon-density.
- ✓ Consequence of $B\bar{B}$ annihilation : $\bar{\Lambda}$ s have higher chance to annihilate with p compared to Λ s that annihilate with \bar{p} causes lowering of low p_T yields while keeping the high p_T yields unchanged, resulting in a hardening of $\bar{\Lambda}$'s p_T -spectra

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$B\bar{B}$ annihilation effect on $\Lambda, \bar{\Lambda} \langle m_T \rangle$



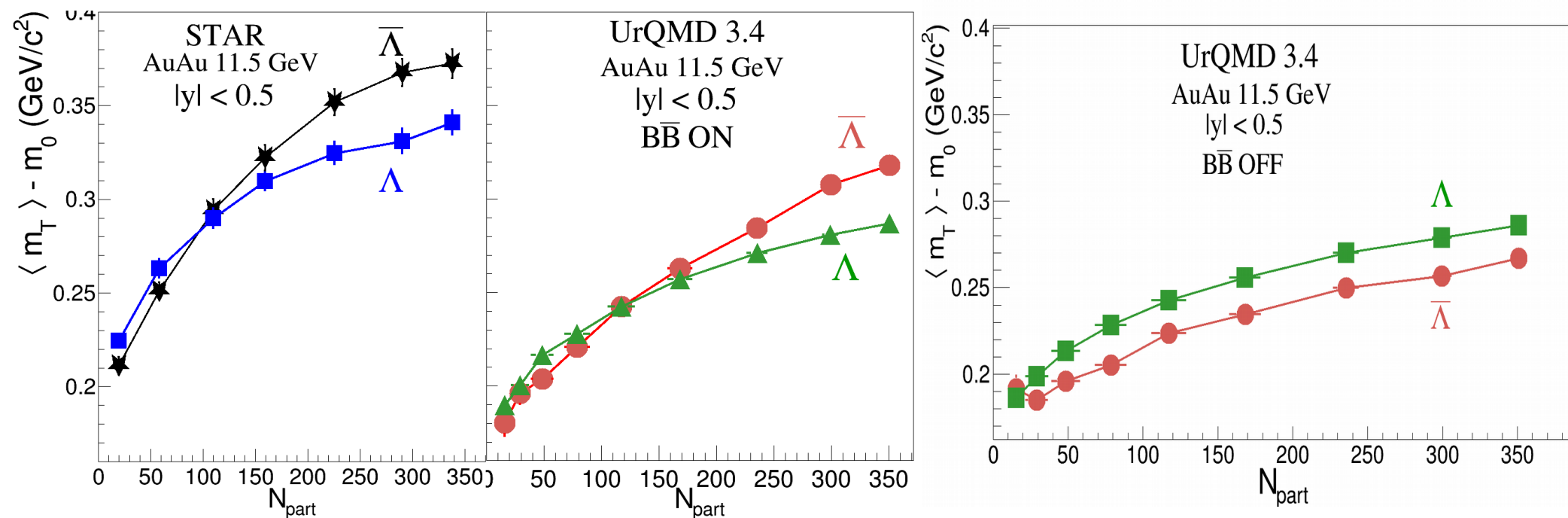
✓ $B\bar{B}$ annihilation off -

$\langle m_T \rangle$ Splitting in opposite trend.

✓ Energy threshold $N + N \rightarrow N + K^+ + \Lambda \sim 700 \text{ MeV}$,
 $N + N \rightarrow \Lambda + \bar{\Lambda} + N + N \sim 2200 \text{ MeV}$

✓ Energy, in excess to threshold energy, available to impart kinetic energy to the production in Λ pair-production is less. Thus, produced $\bar{\Lambda}$ s mostly have smaller transverse momenta.

$\langle m_T \rangle$ vs centrality for Λ and $\bar{\Lambda}$

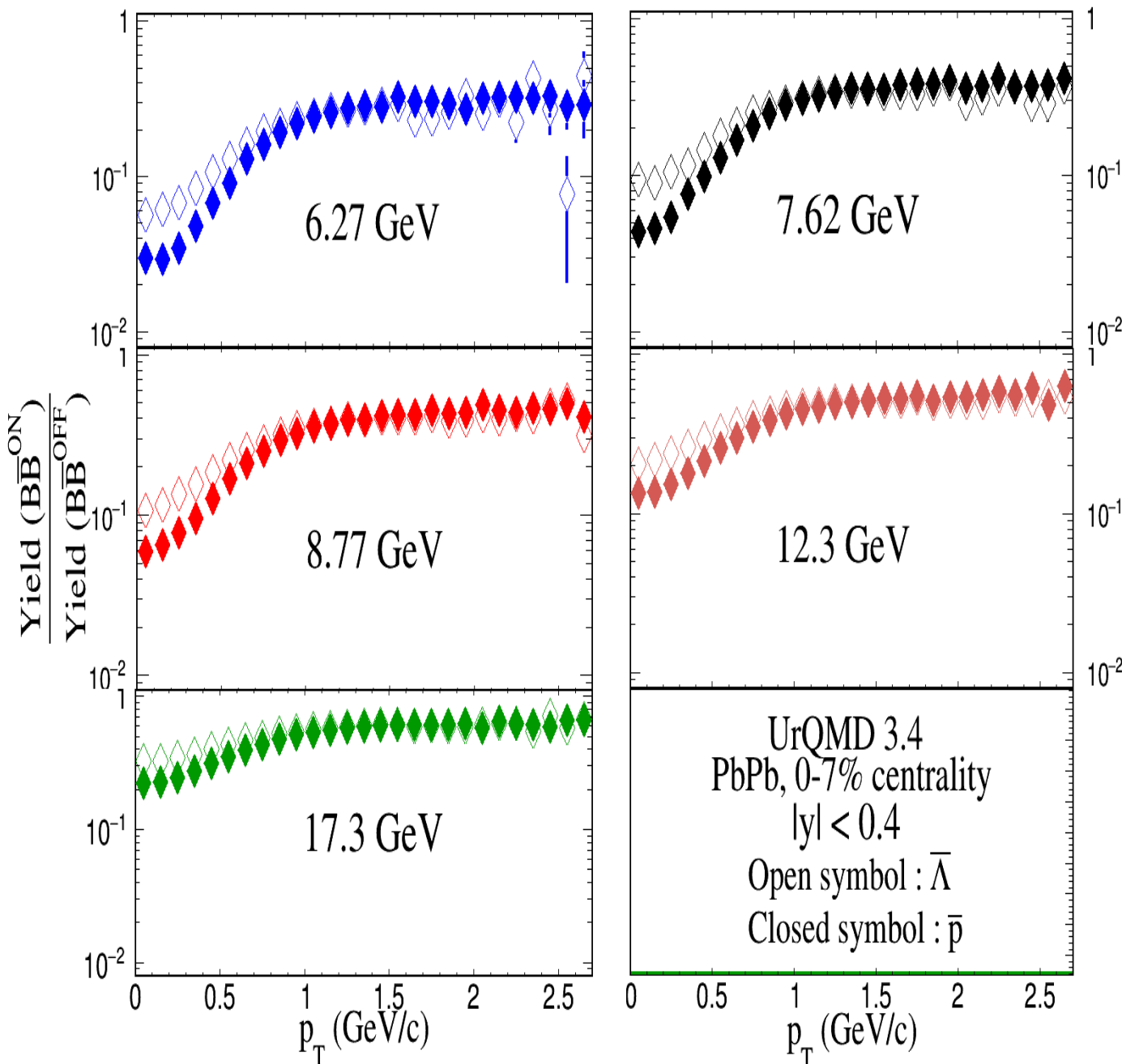


- Shows **interplay between $B\bar{B}$ annihilation and threshold energy**.

$N_{part} < 100$, “threshold effect” is dominant as the baryon density is low. This results in higher values of $\langle m_T \rangle$ for Λ than $\bar{\Lambda}$.

- Switch-over for $N_{part} > 100$, where the trend get reversed i.e, $\langle m_T \rangle$ for $\bar{\Lambda}$ is greater than Λ . This happens because as the system size increases, baryon density increases.
- When $B\bar{B}$ annihilation is OFF in UrQMD, $\langle m_T \rangle$ for Λ is higher than $\bar{\Lambda}$ at all N_{part} .
- Suggests that in central collisions at lower \sqrt{s} effects of $B\bar{B}$ annihilation can not be ignored particularly while considering any phenomenon that involves yields or spectra of baryons and anti-baryons.

p_T dependence of yields of Λ 's & \bar{p} 's



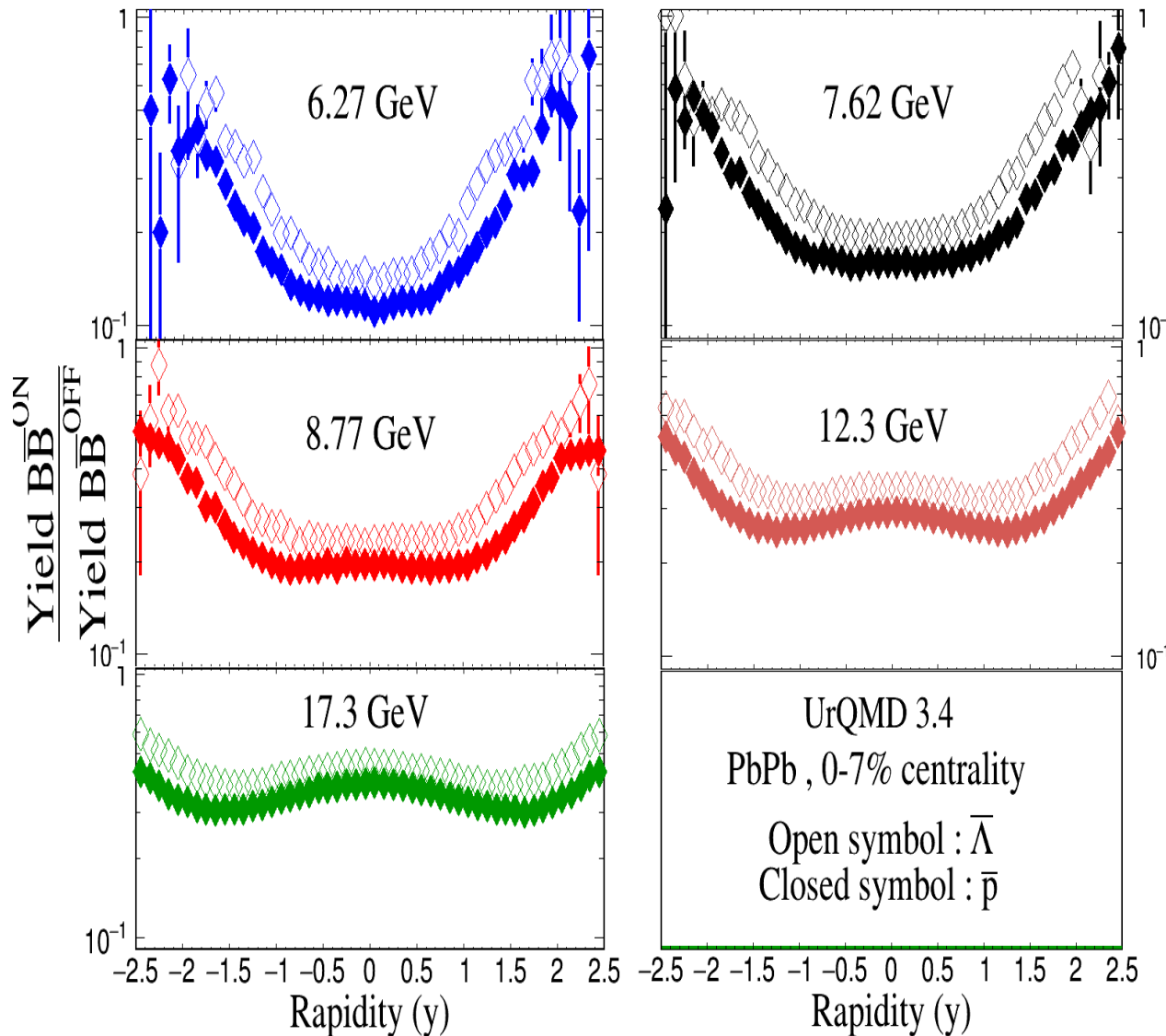
To investigate how $\bar{B}\bar{B}$ annihilation modify yields of Λ and \bar{p} as a function of p_T

Suggests that the effect of annihilation is maximum at low- p_T

\bar{p} is annihilated more than Λ

The magnitude of suppression decreases with increasing \sqrt{s} , implying more annihilation of Λ and \bar{p} at lower \sqrt{s}

Rapidity dependence yields of Λ 's & \bar{p} 's



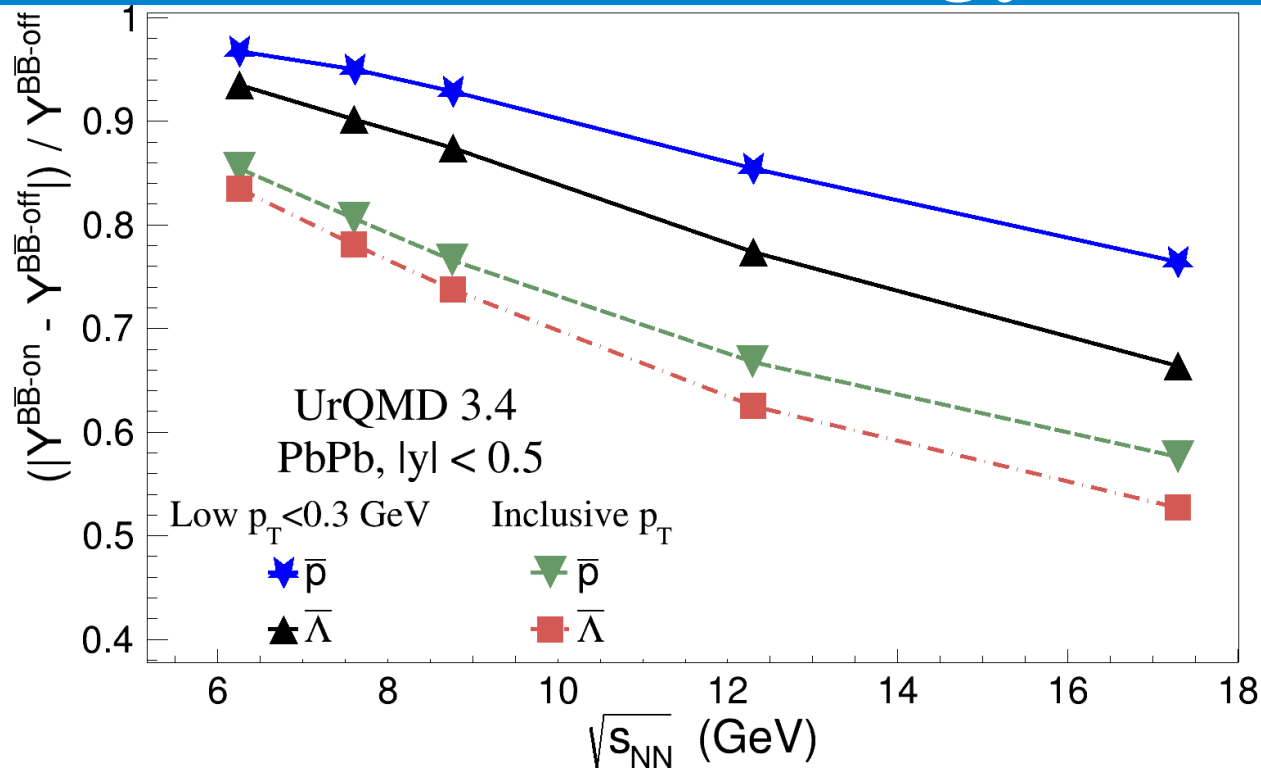
The annihilation effect is most significant at mid-rapidity in low \sqrt{s} for both Λ and \bar{p}

As energy increases, the overall magnitude of suppression start to decrease.

The maximum suppression now occurs at forward rapidity instead of mid-rapidity at higher energy

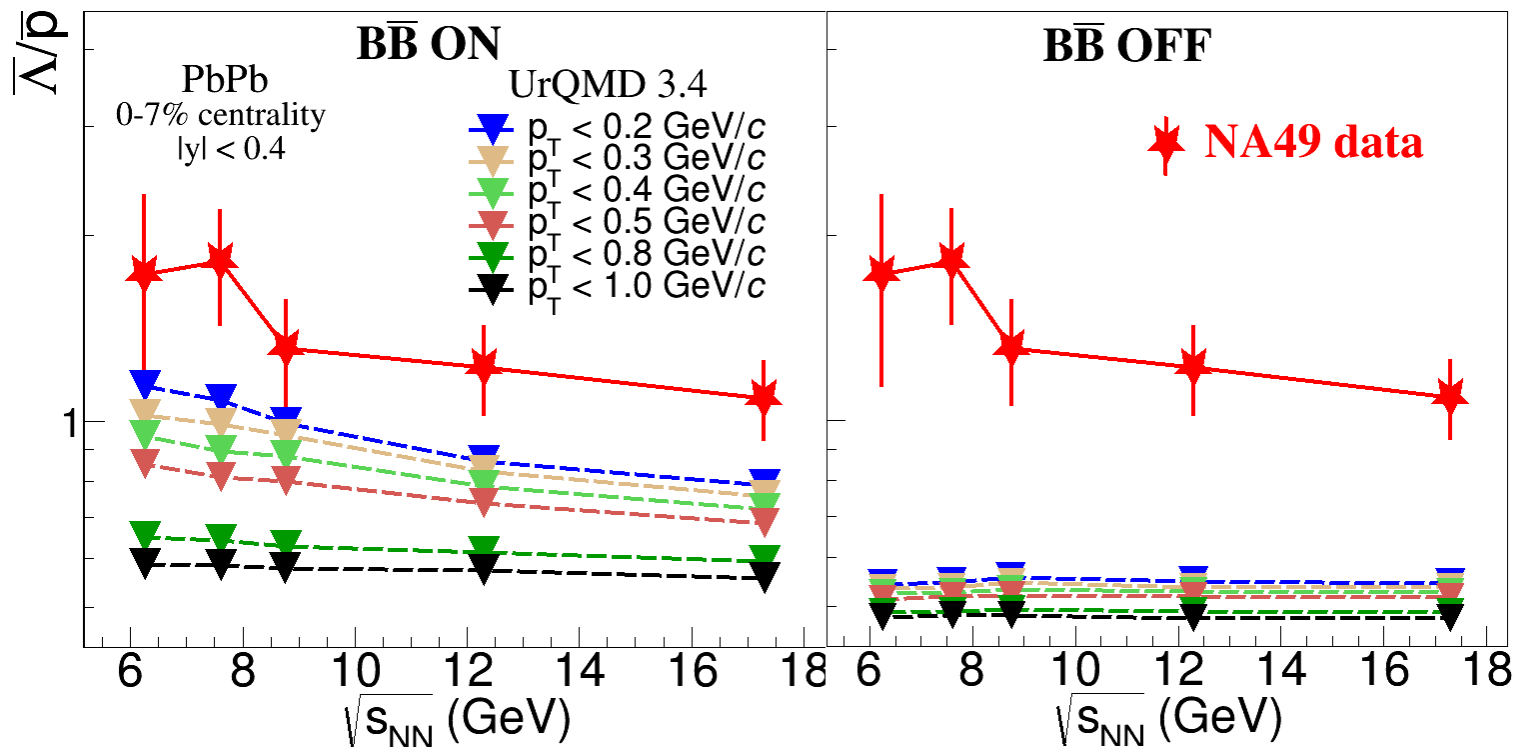
At high energy baryon stopping is less at mid rapidity & high baryon density region gradually move to forward rapidity.

Annihilation fraction as function of beam energy



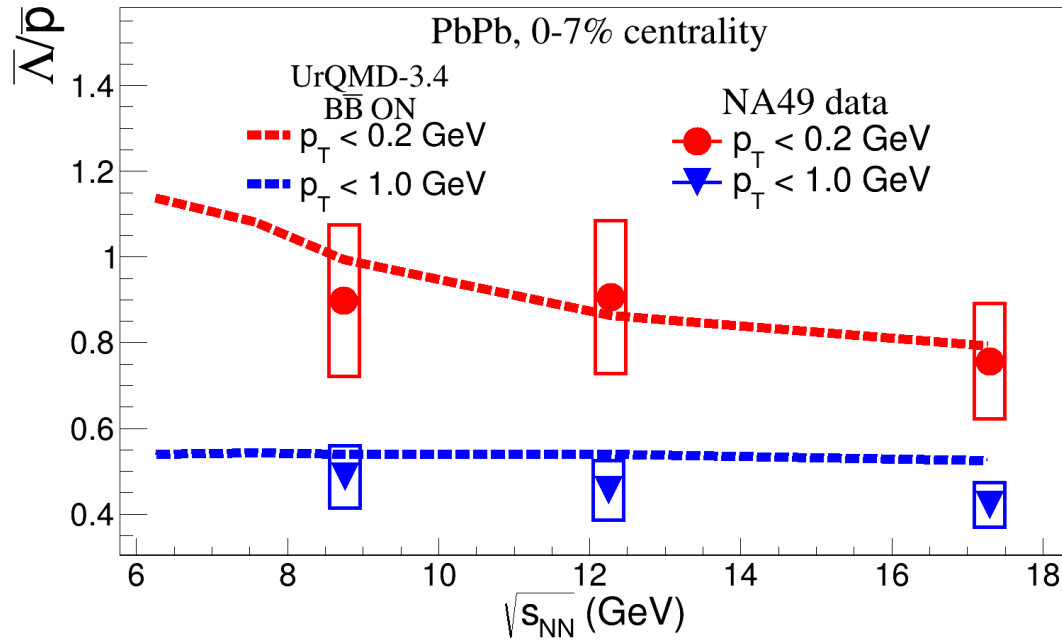
- ✓ Magnitude of annihilation interms of a quantity called **annihilation fraction**.
- ✓ Annihilation is more at low p_T than inclusive p_T .
At low \sqrt{s} , say at 6.27 GeV, $> 95\%$ of the initially produced low- p_T \bar{p} are annihilated whereas, for $\bar{\Lambda}$ the annihilation fraction is about 92 to 94%
- ✓ As the energy increases annihilation fraction drops.
- ✓ Annihilation fraction of \bar{p} is higher than $\bar{\Lambda}$.

The Λ/\bar{p} as a function of collision energy



- ✓ Λ/\bar{p} are sensitive to $B\bar{B}$ annihilation and its impact depend strongly on the kinematic selection.
- ✓ $p_T < 0.2$ GeV/c) Λ/\bar{p} ratios achieve maximum for the lowest collision energy. The trend in UrQMD is qualitatively similar to data
- ✓ Enhancement in Λ/\bar{p} : Annihilation cross-sections of \bar{p} is higher than Λ which results in more suppression of \bar{p} yield compared to Λ , leading to an enhancement.
- ✓ When UrQMD calculations are done without incorporating $B\bar{B}$ -annihilation, irrespective of the choice of p_T range no enhancement can be seen in Λ/\bar{p} ratios.

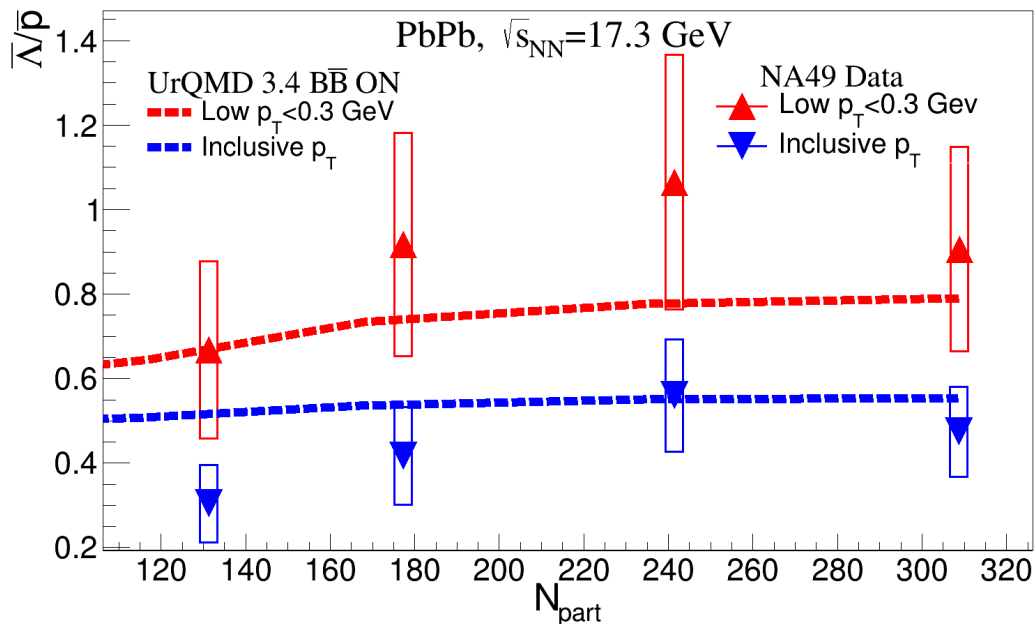
Λ/\bar{p} with collision energy & centrality



◆ we now calculate Λ/\bar{p} ratios from the feed-down corrected yields for both Λ and \bar{p}

◆ After the feed-down correction, Λ/\bar{p} ratios in data and model calculations show very good agreement.

◆ For low- p_T , ratio increases from peripheral to central collision in UrQMD as well as in data



◆ Enhancement in Λ/\bar{p} ratios is only observed at low- p_T but for inclusive p_T ratios are almost flat

◆ Suggests, there is a strong impact of $B\bar{B}$ annihilation on Λ/\bar{p} ratios and the enhancement in Λ/\bar{p} ratios may not be a necessary indication of partonic deconfinement for strangeness enhancement.

Summary

- ✓ $B\bar{B}$ -annihilation has a significant effect on Λ and $\bar{\Lambda}$ hyperon production at high net baryon density or lower energy region.
- ✓ Data model comparison reveals that $B\bar{B}$ annihilation is responsible for the splitting of $\langle m_T \rangle$ as a function of \sqrt{s} for Λ , $\bar{\Lambda}$
- ✓ $\langle m_T \rangle$ with N_{part} shows a interplay between threshold energy and $B\bar{B}$ annihilation effect
- ✓ Λ/\bar{p} ratio enhancement is seen to be well explained by UrQMD with $B\bar{B}$ annihilation at low and inclusive p_T and matches well with feed-down corrected data
- ✓ This investigations suggest Λ/\bar{p} enhancement is not necessarily because of strangeness enhancement due to partonic deconfinement and $B\bar{B}$ -annihilation has a significant role to play.

Thank You

Parametrization of $B\bar{b}$ annihilation in UrQMD

UrQMD use some form parametrization of $B\bar{b}$ annihilation cross section, which are nevertheless data-driven.

$$\sigma_{ann}^{p\bar{p}} = \sigma_0^N \frac{s_0}{s} \left[\frac{A^2 s_0}{(s - s_0)^2 + A^2 s_0} + B \right]. \quad (2)$$

Where $\sigma_0^N = 120$ mb, $s_0 = 4m_N^2$, $A = 50$ MeV and $B = 0.6$ [36].

For annihilation channels that involve a strange-baryon/antibaryon, such as $\bar{\Lambda}p$ or $\Lambda\bar{p}$, an additional correction factor is introduced based on AQM, given by

$$\sigma_{ann}^{B\bar{B}} = (1 - 0.4 \frac{s_B}{3})(1 - 0.4 \frac{s_{\bar{B}}}{3}) \sigma_{ann}^{p\bar{p}} [36], \quad (3)$$

where s_B and $s_{\bar{B}}$ is the strangeness number for baryon and antibaryon, respectively. Thus, annihilation cross-section of $\bar{\Lambda} - p$, ($\sigma_{ann}^{\bar{\Lambda}p}$) is less than annihilation cross-section of $p - \bar{p}$ ($\sigma_{ann}^{p\bar{p}}$). From equation [2] we see the annihilation cross-section (σ_{ann}) has an approximate $\frac{1}{\sqrt{s}}$ dependence with beam energy.

