

Higher-order quark and gluon propagators in thermal medium

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Hard thermal loop Lagrangian (LO)

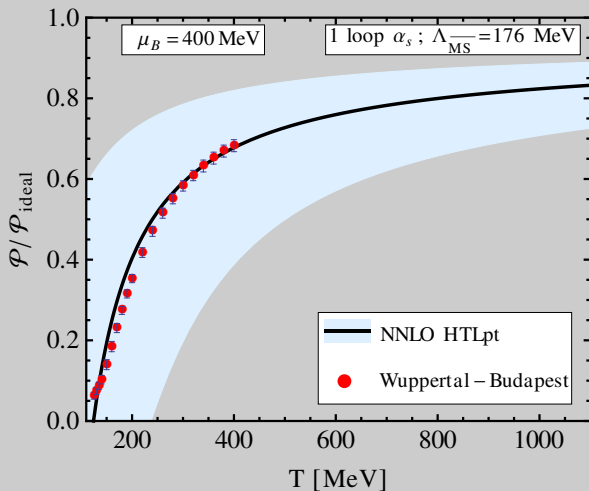
- The HTLpt Lagrangian density (LO) can be written as

$$\mathcal{L} = (\mathcal{L}_{\text{QCD}} + \mathcal{L}_{\text{HTL}})|_{g \rightarrow \sqrt{\delta}g} + \Delta\mathcal{L}_{\text{HTL}},$$

where the HTL improvement term is

$$\mathcal{L}_{\text{HTL}} = (1 - \delta)im_q^2 \bar{\psi} \gamma^\mu \left\langle \frac{y_\mu}{y \cdot D} \right\rangle_{\hat{y}} \psi - \frac{1}{2}(1 - \delta)m_D^2 \text{Tr} \left(F_{\mu\alpha} \left\langle \frac{y^\alpha y_\beta}{(y \cdot D)^2} \right\rangle_{\hat{y}} F^{\mu\beta} \right),$$

- HTL (LO) quarks and gluon propagators are used to study dilepton/photon production rate, quark/gluon damping rate, NNLO QCD thermodynamics and many more.



Curvature of QCD phase transition line

NH et. al., PRC-L(2021)

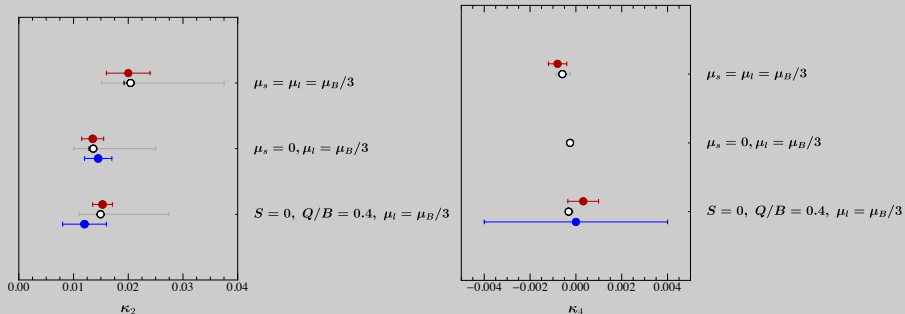


Figure: Filled circles are lattice calculations and black open circles are the NNLO HTLpt predictions

NLO gluon debye mass

- The NLO gluon debye screening mass has been obtained by Rebhan [PRD 48] and Brateen [PRL73] independently using the standard HTL resummation techniques and by calculating the Polyakov loop correlators respectively. The result obtained is

$$\delta m^2 = gm^2 N \sqrt{\frac{6}{2N + N_f}} \frac{1}{2\pi} \left(\ln \frac{2m}{m_{\text{magn}}} - \frac{1}{2} \right) + O(g^2)$$

where $m_{\text{magn}} \sim g^2 T$

Quark NLO

- The quark NLO thermal mass is obtained in ref. [Carrington et. el. PRD 78 (2008)] at zero momentum (for $N_f = 2, N_c = 3$) which is given by

$$M_{\text{QCD}} = \frac{gT}{\sqrt{6}} \left[1 + (1.867 \pm 0.02) \frac{g}{4\pi} \right] + \mathcal{O}(g^3 T),$$

- Since there is no quark damping in the leading order. Damping rate for quarks at NLO results are obtained by E Brateen in ref. [PRD 46]

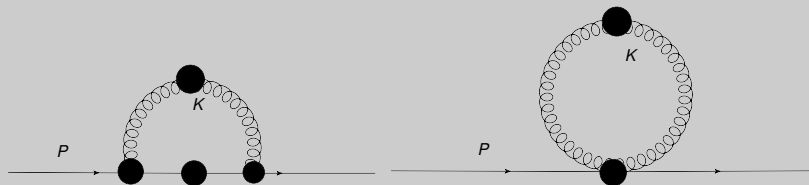
$$\begin{aligned} \gamma_{\pm}(0) &\approx 0.151g^2T \quad (\text{For } N_f = 3) \\ &\approx 0.149g^2T \quad (\text{For } N_f = 2) \end{aligned}$$

- Motivation: To study the quark thermal mass and damping rate at finite quark momentum.

- HTL QED Lagrangian density in NLO

$$\mathcal{L}_{HTL}^{\text{NLO}} = \frac{e^2 v^{3-d}}{2} \int \frac{d^d \mathbf{p}}{(2\pi)^d} \left\{ \frac{N_f(p)}{2p^3} F_{\rho\alpha} \frac{v^\alpha v^\beta}{(v \cdot \partial)^4} \partial^4 F_{\beta\rho} + \right. \\ \left. e^2 (d-1) \Lambda_{(d)}^2 \left(\frac{N_f(p)}{p^3} F_{\rho\alpha} \left[\frac{v^\alpha v^\beta}{(v \cdot \partial)^2} \left(\frac{1}{2} + \frac{\partial_0}{v \cdot \partial} \right) - \frac{n^\alpha v^\beta + v^\alpha n^\beta}{2(v \cdot \partial)^2} \right] F_{\beta\rho} \right. \right. \\ \left. \left. - \frac{1}{2p^2} \frac{dN_f(p)}{dp} F_{\rho\alpha} \frac{v^\alpha v^\beta}{(v \cdot \partial)^2} F_{\beta\rho} \right) \right\},$$

- Full HTL QCD lagrangian in NLO will be an interesting result.



Quark quasi-particle energy

$$\Omega(p) = \Omega^{(0)}(p) + \Omega^{(1)}(p) + \dots$$

Quark self-energy Σ

$$\Sigma(P) = \Sigma_{\text{HTL}}(P) + \Sigma^{(1)}(P) + \dots,$$

where Σ_{HTL} is the lowest-order quark self-energy having gT order, whereas $\Sigma^{(1)}$ the NLO contribution of quark self-energy, with order g^2T .

For a slow-moving quark slow-moving quarks, we can take $p \sim gT$, and we get

$$\Omega_{\pm}^{(1)}(p) = \frac{\Omega_{\pm}^{(0)2}(p) - p^2}{2m_q^2} \Sigma_{\pm}^{(1)}\left(\Omega_{\pm}^{(0)}(p), p\right).$$

We need to determine the NLO quark self-energy to evaluate the above equation.

Quark dispersion relation

- Comparison of quark mass in leading order with NLO in terms of coupling constant g .

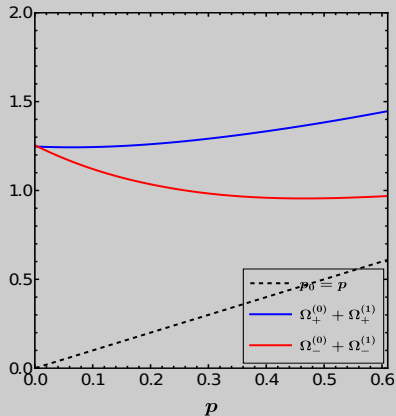
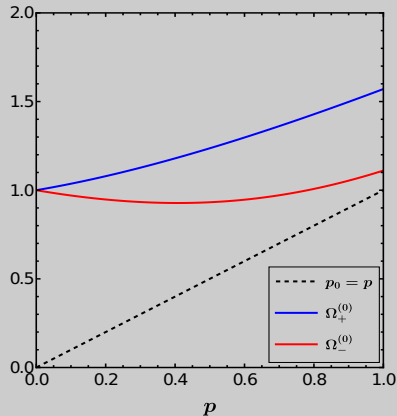


Figure: NLO quark dispersion relation

Quark damping rate

- Damping rate can be obtained from

$$\det(\Phi - \Sigma_{\text{total}}(Q))|_{q_0=M-i\gamma} = 0$$

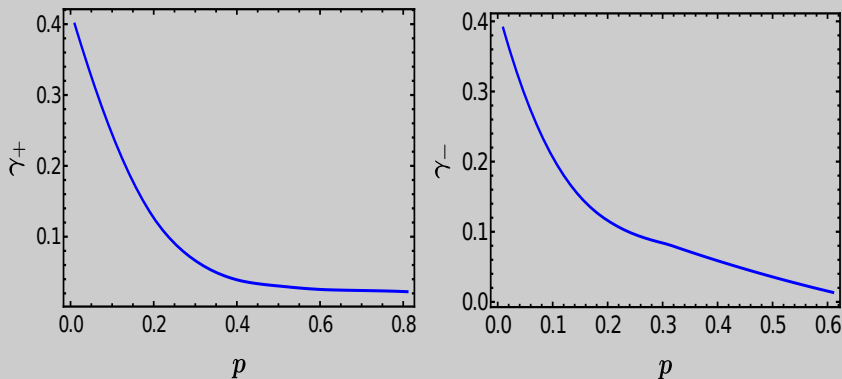


Figure: Quark damping rate

- Velocity variation of quasi-quarks with soft momentum p .

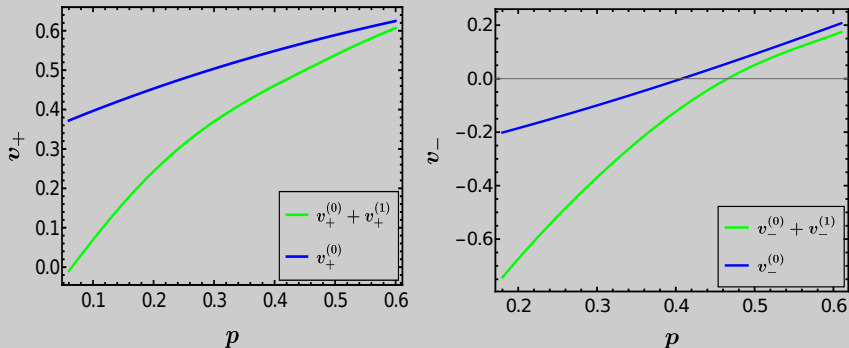


Figure: Velocity variation with respect to soft momentum p for quark mode and plasmino mode.

Thank you for your attention