







Lattice approach to inhomogeneous magnetic fields as probes of QCD thermodynamics

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Motivation	Lattice QCD and magnetic fields	Lattice simulations	Summary & Conclusions	References

- 1. Motivation
- 2. Lattice QCD and magnetic fields
- 3. Lattice simulations
- 4. Summary & Conclusions

Motivation

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PERIPHERAL HICS



 $\sqrt{eB}\sim 0.1~\text{-}~0.5~\text{GeV}$

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Figure 1: Spatial distributions **B** (left) and **E** (right) fields for an impact parameter b = 10 fm. *P* Deng and Huang 2012.

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Figure 1: Spatial distributions **B** (left) and **E** (right) fields for an impact parameter b = 10 fm. *P* Deng and Huang 2012.

Caveats:

- highly non-homogeneous background.
- E leads to sign problem.
- No Minkoswki time evolution in lattice QCD.

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B changes T_c : could the system be in different phases at different x?

Lattice **QCD** and magnetic fields

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$$\langle \ \mathcal{O} \
angle = rac{1}{\mathcal{Z}} \int \mathcal{D} \bar{\psi} \mathcal{D} \psi \mathcal{D} A \ \mathcal{O} e^{-S[\bar{\psi},\psi,A]}$$

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$$\langle \mathcal{O} \rangle = \frac{1}{\mathcal{Z}} \int \mathcal{D}\bar{\psi}\mathcal{D}\psi\mathcal{D}A \ \mathcal{O}e^{-S[\bar{\psi},\psi,A]} \longrightarrow \frac{1}{\mathcal{Z}} \int \mathcal{D}A \det\left[\mathcal{D}(A) + m\right] \mathcal{O}e^{-S_g[A]}$$

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1. Generate samples $\{\mathcal{O}_1, \mathcal{O}_2, ..., \mathcal{O}_N\}$ with a probability $\det [\mathcal{D}(A) + m] e^{-S_g}$ using Monte Carlo steps.

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• magnetic field
$$B \longrightarrow u_{\mu} = e^{iaqA_{\mu}} \in U(1)$$
 (BACKGROUND!)

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UNIFORM MAGNETIC FIELD ON THE LATTICE

$$\mathbf{B} = \mathbf{
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$$A_y = Bx$$

periodic boundary conditions

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periodic boundary conditions

$$u_x = \begin{cases} e^{-iqBL_xy} & \text{if } x = L_x - a\\ 1 & \text{if } x \neq L_x - a \end{cases}$$
$$u_y = e^{iaqBx}$$
$$u_z = 1$$
$$u_t = 1$$

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UNIFORM MAGNETIC FIELD ON THE LATTICE

Flux quantization in a box

$$\mathbf{B} = \mathbf{
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$$A_y = Bx$$

$$qB = \frac{2\pi N_b}{L_x L_y}$$

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INHOMOGENEOUS MAGNETIC FIELD ON THE LATTICE

$$\mathbf{B} = \frac{B}{\cosh\left(\frac{x-L_x/2}{\epsilon}\right)^2} \hat{z}$$
Motivated by HIC scenarios \mathscr{O} Deng and
Huang 2012, \mathscr{O} Cao 2018.

$$qB = \frac{\pi N_b}{L_y \epsilon \tanh\left(\frac{L_x}{2\epsilon}\right)} \qquad N_b \in \mathbb{Z}$$

$$u_x = \begin{cases} e^{-2iqB\epsilon y \tanh\left(\frac{L_x}{2\epsilon}\right)} & \text{if } x = L_x - a\\ 1 & \text{if } x \neq L_x - a\\ 1 & \text{if } x \neq L_x - a \end{cases}$$

$$u_y = e^{iaqB\epsilon [\tanh\left(\frac{x-L_x/2}{\epsilon}\right) + \tanh\left(\frac{L_x}{2\epsilon}\right)]}$$

$$u_z = u_t = 1$$

Lattice simulations

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• $N_f = 2 + 1$ improved staggered fermions with physical masses;

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- $N_f = 2 + 1$ improved staggered fermions with physical masses;
- Lattices: $16^3 \times 6 \quad 24^3 \times 8 \quad 28^3 \times 10 \quad 36^3 \times 12 \quad \longrightarrow$ continuum limit (lattice spacing $\rightarrow 0, V = \text{const.}$);

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$$\mathbf{B} = \frac{B}{\cosh\left(\frac{x - L_x/2}{\epsilon}\right)^2} \hat{z} \qquad eB = \frac{3\pi N_b}{L_y \epsilon \tanh\left(\frac{L_x}{2\epsilon}\right)} \qquad \epsilon \approx 0.6 \text{ fm}$$

strength $0~{\rm GeV} \leq \sqrt{eB} \leq 1.2~{\rm GeV} \longrightarrow$ magnetars, HIC and early universe.

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• Temperature $\longrightarrow 68 \text{ MeV} \le T \le 300 \text{ MeV}$ (crossover at $T_c \sim 155 \text{ MeV}$).

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- Temperature $\longrightarrow 68 \text{ MeV} \le T \le 300 \text{ MeV}$ (crossover at $T_c \sim 155 \text{ MeV}$).
- Quantities: $\Sigma(x)_{T,B} = P(x)_{T,B} = \mathbf{J}(x)_{T,B}$.

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Enchancement of the condensate: magnetic catalysis.

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$$C(x,x') = \frac{1}{m_\pi^3} \left\langle \ \bar{\psi}\psi(x)P(x') \ \right\rangle_c$$

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 $P(x)_{T,B}$ is broader than the chiral condensate.

$$C(x,x') = \frac{1}{m_{\pi}^3} \left\langle \ \bar{\psi}\psi(x)P(x') \ \right\rangle_c$$



Interaction with *P* causes the dips! (Local inverse magnetic catalysis)

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Electric currents: $\mathbf{J} \sim \mathbf{ abla} imes \mathbf{B}$

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ELECTRIC CURRENTS: $\mathbf{J} \sim \boldsymbol{ abla} imes \mathbf{B}$



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ELECTRIC CURRENTS: $\mathbf{J} \sim \mathbf{ abla} \times \mathbf{B}$



Figure 4: Lattice electric currents for RHIC-like ($\sqrt{eB} = 0.1 \text{ GeV}$) and LHC-like ($\sqrt{eB} = 0.5 \text{ GeV}$) magnetic fields, respectively.

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$$\frac{1}{\mu_0}\mathbf{B} = \mathbf{H} + \mathbf{M} \qquad \mathbf{J}_m = \mathbf{\nabla} \times \mathbf{M}$$

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$$rac{1}{\mu_0} \mathbf{B} = \mathbf{H} + \mathbf{M} \qquad \mathbf{J}_m = \mathbf{
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Linear response term:

 $\mathbf{M} \approx \chi_m \mathbf{H}$

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Linear response term:

$$\begin{split} \mathbf{M} &\approx \chi_m \mathbf{H} \\ \frac{\chi_m}{1 + \chi_m} \mathbf{\nabla} \times \mathbf{B} = \mathbf{J}_m \end{split}$$

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The divergence is independent of T: $\chi_m^r(T) \equiv \chi_m(T) - \chi_m(0)$

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Great agreement other predictions / Bali, Endrődi, and Piemonte 2020

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$$\left. \frac{\partial J^5}{\partial \mu} \right|_{\mu=0} = C_{\rm CSE} \ eB$$

See the talk by C. Garnacho (Tuesday

at 17:40).



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$$\left. \frac{\partial J^5}{\partial \mu} \right|_{\mu=0} = C_{\rm CSE} \ eB$$

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 $16^{3} \times 6$

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What about CME? Work in progress!

Summary & Conclusions

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• Richer scenario with B(x) (dips, electric currents, etc.)

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- Richer scenario with B(x) (dips, electric currents, etc.)
- Observables change significantly due to B(x)

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- Richer scenario with B(x) (dips, electric currents, etc.)
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- New method to extract χ_m using $\mathbf{J}(x)$ and Maxwell's eqs.
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SUMMARY & CONCLUSIONS

- Richer scenario with B(x) (dips, electric currents, etc.)
- Observables change significantly due to B(x)
- New method to extract χ_m using J(x) and Maxwell's eqs.
- CSE does not feel the inhomogeneity of B (CME?)

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SUMMARY & CONCLUSIONS

- Richer scenario with B(x) (dips, electric currents, etc.)
- Observables change significantly due to B(x)
- New method to extract χ_m using J(x) and Maxwell's eqs.
- CSE does not feel the inhomogeneity of *B* (CME?)
- *B*(*x*) is important to capture the correct physics in peripheral HICs (applications to QCD models, hydrodynamics, etc.)

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धन्यवाद!

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BIBLIOGRAPHY I

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