Aspects of chiral transition in a Hadron Resonance Gas Model

Deeptak Biswas

The Institute of Mathematical Sciences

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In collaboration with Peter Petreczky, Sayantan Sharma

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Introduction

Chiral Condensate in the Hadron Resonance Gas model Renormalized definitions of chiral condensate

Results

HRG results of renormalized chiral condensate vs. LQCD Curvature of the chiral crossover line from HRG model

Summary and outlook

Introduction

Motivation

- ➤ In QCD with 2 massless quarks, $SU(2)_{\nu} \times SU(2)_{A} \times U(1)_{\nu} \rightarrow \text{Exact symmetry}$
- ➤ With physical mass, SU(2)_v × SU(2)_A × U(1)_v → Approximate (good)

> This symmetry is spontaneously broken to $SU(2)_{\nu} \times U(1)_{\nu}$

- For 2 flavors at non-zero mass, the chiral symmetry is restored via analytic crossover at $T_c = 156.5(1.5)$ MeV. [HotQCD 2018]
- How far can we estimate T_c and pseudo-critical line in Hadron Resonance Gas(HRG) model?



Chiral Condensate in the Hadron Resonance Gas model

The earliest estimation of pseudo-critical temperature, done within the NNLO chiral perturbation theory(χ_{PT}) gave $T_c = 250$ MeV.

[P. Gerber, H. Leutwyler 1989]

- Lowered to about 190 MeV with inclusion of heavier hadrons.
- \blacksquare Recent studies within the HRG have found a higher $T_c \sim 170$ MeV.

[J. Jankowski et al. 2013, A. N Tawfik, N. Magdy 2015]

Renormalized chiral condensate

We can define the renormalized chiral condensate from the pressure as,

$$-m_{s}\left[\langle\bar{\psi}\psi\rangle_{I,T}-\langle\bar{\psi}\psi\rangle_{I,0}\right]=-m_{s}\frac{\partial P}{\partial m_{I}}$$

The normalization is not unique [BMW 2010],

$$\langle \bar{\psi}\psi \rangle_{R} = -\frac{m_{I}}{m_{\pi}^{4}} \left[\langle \bar{\psi}\psi \rangle_{I,T} - \langle \bar{\psi}\psi \rangle_{I,0} \right]$$

A natural choice for dimensionless condensate [HotQCD 2012],

$$\Delta_{R}^{I} = d + m_{s} r_{1}^{4} \left[\langle \bar{\psi}\psi \rangle_{I,T} - \langle \bar{\psi}\psi \rangle_{I,0} \right]$$

■ Using low energy constant of $SU(2) \chi_{PT}$, $\Sigma^{1/3} = 272(5)$ MeV, $m_s = 92.2(1.0)$ MeV, and $r_1 = 0.3106$ fm, one gets d = 0.022791. [FLAG 2022],

The renormalized chiral condensate,

$$m_s \frac{\partial P}{\partial m_l} = -\frac{m_s}{m_l} \sum_{\alpha} \frac{g_{\alpha}}{2\pi^2} \int_0^{\infty} dp \ p^2 \ n_{\alpha} \ (E_{\alpha}) \frac{1}{2E_{\alpha}} m_l \frac{\partial M_{\alpha}^2}{\partial m_l}.$$

The non-trivial ingredient is $\frac{\partial M_{\alpha}^2}{\partial m_l}$.

From $SU(2) \chi_{PT}$,

$$M_{\pi}^2 = M^2 \left[1 - \frac{1}{2} \zeta \ \bar{l}_3 + \mathcal{O}(\zeta^2) \right] \quad , \quad \zeta = \frac{M^2}{16\pi^2 F_{\pi}^2}$$

- Kaon properties are predicted well from 2+1 χ_{PT} [RBC 2014, Durr 2015] $M_K^2 = B_K(m_s)m_s \left[1 + \frac{\lambda_1(m_s) + \lambda_2(m_s)}{F^2}M^2\right]$ $M^2 = 2Bm_l, B = \Sigma/F^2$
- From LQCD the pion mass is consistent with LO result $M_{\pi}^2 \approx 2Bm_I$. [RQCD Bali et al. 2016].

Sigma terms for Heavier hadrons

$$\sigma_{\alpha} = m_{l} \frac{\partial M_{\alpha}}{\partial m_{l}}|_{m_{l} = m_{l}^{phys}} = m_{l} \langle \alpha | \bar{u}u + \bar{d}d | \alpha \rangle = M_{\pi}^{2} \frac{\partial M_{\alpha}}{\partial M_{\pi}^{2}}|_{M_{\pi} = M_{\pi}^{phys}}.$$

Ν	٨	Σ	Ξ
44(3)(3)	31(1)(2)	25(1)(1)	15(1)(1)
Δ	Σ^*	Ξ*	Ω^{-}
29(9)(3)	18(6)(2)	10(3)(2)	5(1)(1)

The sigma terms of ground state baryons have been only recently calculated with precision. [Copeland et al. 2021].

New development from our work:

- We have done extensive compilation of the LQCD results to find $M_{\pi}^2 \frac{\partial M_{\alpha}}{\partial M_{\pi}^2}$ at a constant m_s , set at the physical value.
- For the first time, σ terms for η , $\rho(770)$, $K^*(892)$, and η' have been calculated from LQCD data.

 $\left[\mathsf{RQCD} \text{ Bali et al. 2016, D. Guo et al. 2016, } \mathsf{RQCD} \text{ Bali et al. 2021} \right]$.

- We have assigned sigma terms for all meson resonances,
 - Iso-vector mesons $\rightarrow \sigma_{\rho(770)}$
 - Open strange mesons $\rightarrow \sigma_{K^*(892)}$.
 - \bullet Iso-scalar mesons \rightarrow corresponding ground states mesons σ terms.
- $f_0(500)$ is not considered as cancellation by the repulsive interactions [Broniowski et al. 2015].

Sigma terms for Baryon resonances: Nucleons

- It is difficult to measure baryon resonances in LQCD as they are close to the scattering state and resonances.
- For the excited N state, the fit to 2+1 flavor LQCD data gives $\sigma = 68(27)$ MeV.
- Within large errors is consistent with the sigma term of its ground state.



- > We have considered the σ terms for all resonances (even for strange baryons) to be same as the ground state.
- To reliably account for large uncertainty in σ of high mass resonances, we have taken the relative errors in the σ-terms of excited states to be 50%.
- However such large uncertainty contributes to only 10% of the total error in the renormalized chiral condensate as the dominant contribution comes from ground state pseudo-scalar mesons.

Results

Chiral condensate: LQCD vs. HRG model



HRG model calculations are consistent with LQCD continuum estimates till T \sim 140MeV.

$$\Delta_{R}^{\prime} = d + m_{s} r_{1}^{4} \left[\langle \bar{\psi}\psi \rangle_{I,T} - \langle \bar{\psi}\psi \rangle_{I,0} \right]$$



- On the lattice Δ_R^l goes to half of its low-temperature value at T_c .
- We use this fact to estimate T_c from our HRG model calculations.

- > Our improved HRG calculation gives $T_c = 161.2 \pm 1.7$ MeV at $\mu_B = 0$.
- > Lattice QCD results on T_c in the continuum limit, $T_c = 156.5 \pm 1.5$ MeV [HotQCD 2018, BMW 2020]

Curvature of the pseudo-critical line

We extract κ_2 and κ_4 by fitting $T_c(\mu_B)$ for $0 < \mu_B/T_c(\mu_B = 0) < 1$.

- Our estimation $\kappa_2 = 0.0203(7)$.
- $\kappa_4 = -3(2) \times 10^{-4}$ is quite noisy.
- Highlight that our results are in very good agreement with LQCD estimates of κ₂ = 0.012(4) [HotQCD 2018] and 0.0153(18) [BMW 2020], κ₄ = 0.



Transition in the chiral limit

➡ 3-loop χ_{PT} for pions + hadrons, gave a $T_c^0 \sim 170$ MeV.

[P. Gerber, H. Leutwyler 1989]

- ⇒ 3-loop χ_{PT} + our improved HRG estimates lower it to 162 MeV.
- LQCD predicts $T_c = 132^{+3}_{-6}$ MeV [HotQCD 2019].
- Need to go beyond 3-loop *χ_{PT}* and include temperature dependent width of the crit- ical *f*₀ mode to improve the agreement



Summary

- We have studied chiral observables for physical hadrons within the HRG model.
- ✓ For the first time, precise values of σ terms for ρ, η, K^* , isoscalar mesons and ground state baryons have been included.
- This has successfully improved the T_c from HRG model, bringing it closer to the LQCD estimates.
- ✓ Curvature coefficients κ_2 , κ_4 are very close to lattice results than previous estimates.