Impact of Baryon anti-Baryon annihilation on apparent strangeness enhancement in /\[v]p at SPS energy

Ekata Nandy

Variable Energy Cyclotron Centre (VECC) Experimental High Energy Physics & Application Group

Outline

✓ Strangeness enhancement

✓ Measures of strangeness enhancement.

✓ Importance of anti-lambda to anti-proton ratio (Λ/\bar{p})

✓ Effects of hadronic interactions importantly baryon-anti- baryon annihilation on Λ , $\Lambda \Lambda/\bar{p}$.

Baryon-anti- baryon $(B\overline{B})$ annihilation effect on

- $\odot < m_{_{
 m T}} > {
 m of} \ \Lambda, \ \Lambda$
- $= \langle m_{T} \rangle$ with Centrality of Λ , Λ
- **•** Yield Λ , \bar{p} w.r.to p_{T} and rapidity
- Annihilation fraction with beam energy
- I√p̄ Ratio

✓ Summary

Ref: Ekata Nandy and Subhasis Chattopadhyay, European Physical Journal A, 58 (2022) 10, 199

Phase diagram of nuclear matter

- ✓ Under extreme conditions of temperature or pressure normal nuclear matter (hadronic phase) is likely to undergo a deconfinement phase transition to a quark gluon plasma phase.
- ✓ Effort to locate the phase boundary at different regions using different experiments.



LHC – Large Hadron Collider (E_{cm} = 2.76TeV – 5.02 TeV)

RHIC BES – Relativistic Heavy Ion Collider Beam Energy Scan (E_{cm} = 7.7 GeV – 200 GeV)

FAIR – Facility for Antiproton & Ion Research (E_{cm} = 3 GeV – 9 GeV)

Signatures of QGP

There is no unique signal that will identify QGP. Different signatures are used to search for QGP.

- J/Ψ suppression
- Strangeness enhancement
- Jet quenching
- Dilepton production

Strangeness production

J. Rafelski and B. Müller first predicted Strangeness enhancement as a signature of deconfinement As there is no initial valence strange quark, it produces from the reactions only.



Hadronic channel

 $N + N \rightarrow N + K + \Lambda,$ $N + N \rightarrow \Lambda + \Lambda + N + N$ $\pi^{+} + N \rightarrow \Lambda + K^{+}$

Why do we expect strangeness enhancement at low energy? (Fermi Energy and Pauli Blocking)

- Because of higher abundance light quarks (u,d) in the medium they fill up the available low energy levels upto the fermi energy. Thus to produce a uu pair, required energy = fermi energy + $2m_{\mu}$
- Thus it is energetically favourable to produce ss pairs that require a threshold energy just double the mass of strange quark only.

Measures of strangeness enhancement

- A large enhancement in strange hadrons production relative to pp interaction was first reported at CERN-SPS.
- **Enhanceent factor** $\frac{dN(Pb+Pb)}{dy}\Big|_{y=0}\Big/\frac{dN(p+p)}{dy}\Big|_{y=0}\Big]$
- Interesting structures were observed in the strangeto-non-strange particle ratios.
- Non-monotonic variation of k/π as a function of collision energy was first observed at SPS.
- Similar behaviour was also observed in the baryon sector (\hbar/\bar{p}), although with large uncertainty.
- Such non-monotonic variation is often attributed to the onset of partonic deconfinement.
- However hadronic interactions such as BB annihilation has a dominant role to play in strangeness baryon production & enhancement in the baryon sector.



Motivation of work

- Understand the contribution of hadronic and partonic sources to the measures of strangeness enhancement.
- Final yields of baryons are highly sensitive to hadronic interactions at later stages of the collisions mainly from the baryon-anti baryon annihilation.
- Such annihilation processes have significant effect on the final yields in a baryon rich environment (CBM & SPS energies). So depending on the different annihilation cross-section of \bar{p} and Λ , this ratio (Λ/\bar{p}) may enhance.
- This study further aims to address whether the enhancement in the ratio (Λ/\bar{p}) can be explained from the consequence of hadronic interactions alone ?

Why Λ/\bar{p} ?

Since anti-particles comprise of quarks produced in the reactions only, they are regarded as a cleaner channel to probe strangeness enhancement

Details of model simulation

UrQMD (Ultra Relativistic Quantum Molecular Dynamics)

- Hadronic model, describe the phenomenology of particle production in pp, p+A and A+A collisions over a broad energy range
- The underlying degrees of freedom in UrQMD are hadrons and strings.
- Here an individual particle propagates on a straight line until the relative distance between two particles is smaller than the total interaction cross-section between two particles.
- The particle production dynamics is either governed by the decays of baryon or meson resonances or via a string excitation and fragmentation.

System : Pb+Pb/Au+Au With & w/o incorporating baryon -anti baryon annihilation Energy : 6.27 GeV, 7.62 GeV, 8.77 GeV, 12.3 GeV, 17.3 GeV Centrality = 0-7% Observables : Λ , \bar{p}

$B\bar{B}$ annihilation effect on Λ , $\Lambda < m_{T} >$



• Average transverse mass \equiv effective temperature. (T_{eff}) . T_{eff} = T_{th} + 1/2 m β^2

- ✓ Although Λ and Λ have same mass, their $\langle m_T \rangle$ are different in magnitude and the difference is seen to increase with the decrease in beam energy or increase in net baryon-density.
- ✓ Consequence of BB annihilation : As have higher chance to annihilate with p compared to As that annihilate with p̄ causes lowering of low p_T yields while keeping the high p_T yields unchanged, resulting in a hardening of A's p_T -spectra
 Ref : Ekata Nandy and Subhasis Chattopadhyay, ⁸
 European Physical Journal A, 58 (2022) 10, 199

$B\bar{B}$ annihilation effect on Λ , $\Lambda < m_{T} >$



• BB annihilation off -

 $< m_{_{\rm T}}>$ Spiltting in opposite trend.

- ✓ Energy threshold N + N → N+ K⁺ + $\Lambda \sim$ 700MeV, N + N → $\Lambda + \Lambda + N + N \sim$ 2200 MeV
- ✓ Energy, in excess to threshold energy, available to impart kinetic energy to the production in ∧ pairproduction is less. Thus, produced ∧s mostly have smaller transverse momenta.

$< m_{\tau} > vs$ centrality for Λ and Λ



✓ Shows interplay between $B\bar{B}$ annihilation and threshold energy. N_{part} <100, "threshold effect" is dominant as the baryon density is low. This results in higher values of $<m_{T}>$ for Λ than Λ .

- ✓ Switch-over for $N_{part} > 100$, where the trend get reversed i.e, $<m_T >$ for Λ is greater than Λ . This happens because as the system size increases, baryon density increases.
- ✓ When $B\bar{B}$ annihilation is OFF in UrQMD, $< m_T > \text{ for } \Lambda$ is higher than Λ at all $N_{_{\text{part}}}$.
- ✓ Suggests that in central collisions at lower √s effects of BB annihilation can not be ignored particularly while considering any phenomenon that involves yields or spectra of baryons and anti-baryons.

p_T dependence of yields of Λ 's & \bar{p} 's



Rapidity dependence yields of A's & p's



The annihilation effect is most significant at mid-rapidity in low \sqrt{s} for both Λ and \bar{p}

As energy increases, the overall magnitude of suppression start to decrease.

The maximum suppression now occurs at forward rapidity instead of mid-rapidity at higher energy

At high energy baryon stopping is less at mid rapidity & high baryon density region gradually move to forward rapidity.

Annihilation fraction as function of beam energy



✓ Magnitude of annihilation interms of a quantity called annihilation fraction.

- \checkmark As the energy increases annihilation fraction drops.
- It is seen that annihilation fraction of Λ drops faster than \bar{p} as a function \sqrt{s} .

The Λ/\bar{p} as a function of collision energy



✓ Λ/\bar{p} are sensitive to $B\bar{B}$ annihilation and its impact depend strongly on the kinematic selection.

- ✓ p_T < 0.2 GeV/c) Λ/\bar{p} ratios achieve maximum for the lowest collision energy. The trend in UrQMD is qualitatively similar to data
- ✓ Enhancement in Λ/\bar{p} : Annihilation cross-sections of Λ is less than \bar{p} which results in more suppression of \bar{p} yield compared to Λ , leading to an enhancement.
- ✓ Interesting to note when UrQMD calculations are done without incorporating $B\bar{B}$ -annihilation, irrespective of the choice of p_{T} range no enhancement can be seen in Λ/\bar{p} ratios.

Λ/\bar{p} with collision energy & centrality



- we now calculate Λ/p̄ ratios from the feed-down corrected yields for both Λ and p̄
- After the feed-down correction, \(\Lambda \/\p\) ratios in data and model calculations show very good agreement.
- Enhancement in \(\Lambda / \bar{p}\) ratios is only observed at low- p_T but for inclusive p_T ratios are almost flat
- For low- p_T, ratio increases from peripheral to central collision in UrQMD as well as in data
- Suggests, there is a strong impact of BB
 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

 BB

Summary

- **V** $B\bar{B}$ -annihilation has a significant effect on Λ and Λ hyperon production at SPS energies.
- ♦ Data model comparison reveals that $B\bar{B}$ annihilation is responsible for the splitting of $< m_T >$ as a function of \sqrt{s}
- ▶ N_{part} dependence of $\langle m_T \rangle$ shows a competing effect threshold energy and $B\bar{B}$ annihilation causes a switch-over in the splitting of $\langle m_T \rangle$
- ← Enhancement in Λ/\bar{p} is seen to be well explained by UrQMD with $B\bar{B}$ annihilation at low and inclusive p_{τ} after considering feed-down correction
- ✓ This investigations suggest Λ/\bar{p} enhancement is not necessarily because of strangeness enhancement due to partonic deconfinement and BB-annihilation has a significant role to play.

Thank You

Parametrization of Bbbar annihilation in UrQMD

UrQMD use some form parametrization of Bbar annihilation cross section, which are nevertheless data-driven.

$$\sigma_{ann}^{p\bar{p}} = \sigma_0^{N} \frac{s_0}{s} \left[\frac{A^2 s_0}{(s-s_0)^2 + A^2 s_0} + B \right].$$
(2)

Where $\sigma_0^N = 120 \text{ mb}$, $s_0 = 4m_N^2$, A = 50 MeVand B = 0.6 [36].

For annihilation channels that involve a strange-baryon/antibaryon, such as $\bar{\Lambda}p$ or $\Lambda\bar{p}$, an additional correction factor is introduced based on AQM, given by

$$\sigma_{ann}^{\mathrm{B}\bar{\mathrm{B}}} = (1 - 0.4 \frac{s_B}{3})(1 - 0.4 \frac{s_{\bar{B}}}{3})\sigma_{ann}^{\mathrm{p}\bar{\mathrm{p}}}[36], \quad (3)$$

where s_B and $s_{\bar{B}}$ is the strangeness number for baryon and antibaryon, respectively. Thus, annihilation cross-section of $\bar{\Lambda} - p$, $(\sigma_{ann}^{\bar{\Lambda}p})$ is less than annihilation cross-section of p- \bar{p} ($\sigma_{ann}^{p\bar{p}}$). From equation [2] we see the annihilation cross-section (σ_{ann}) has an approximate $\frac{1}{\sqrt{s}}$ dependence with beam energy.

