

# Virtual photon polarization and dilepton anisotropy in relativistic heavy-ion collisions

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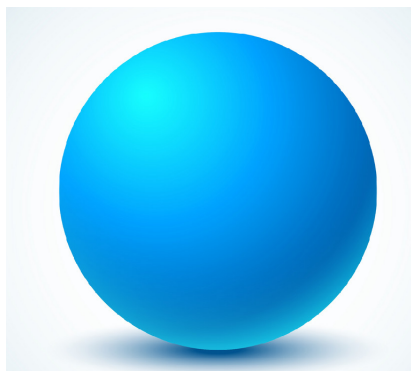
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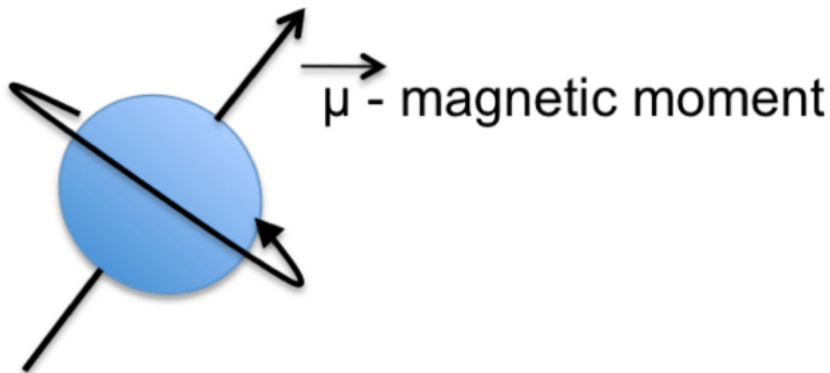
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# Decay of scalar particles



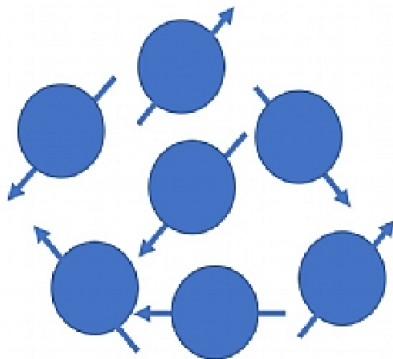
No anisotropy in the rest frame: isotropic decay products.

## Decay of particles with spin



Preferred direction due to spin: anisotropic decay products  
Leads to polarization observables.

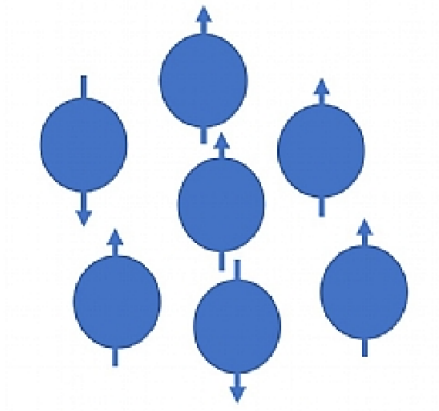
## Several random decays



Expectation: Averaging over random decays should lead to isotropic decay products.

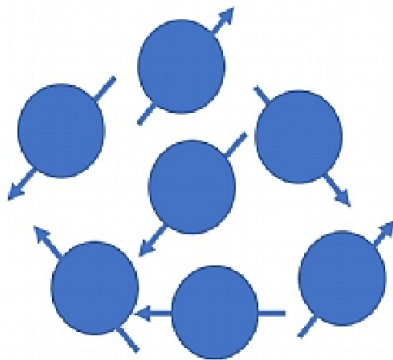
Averaging with thermal distribution.

# Decay of spin polarized particles



Averaging over decay of polarized particles should lead to anisotropic decay products.

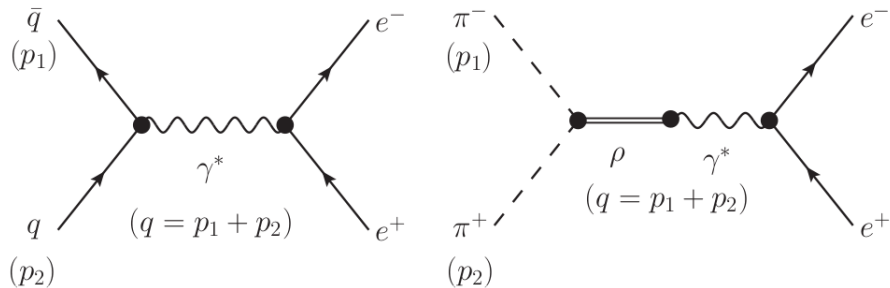
## Several random decays



Averaging over random decays also leads to anisotropic decay products. [PLB 782 \(2018\) 395-400.](#)

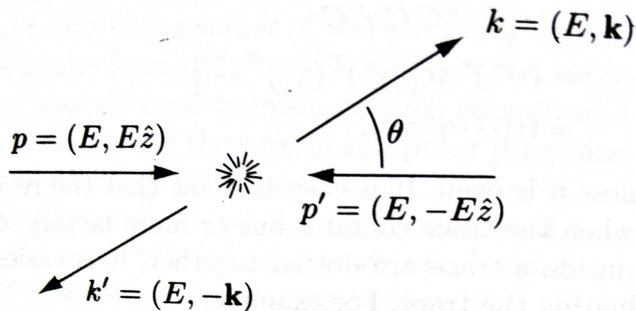
Averaging with thermal distribution.

# Processes under consideration



Feynman diagrams for the Drell-Yan and pion annihilation.

$$e^+ e^- \rightarrow \mu^+ \mu^-$$



$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E_{\text{cm}}^2} (1 + \cos^2 \theta)$$



# Dilepton production rate

- Consider the general process  $X_1 X_2 \rightarrow \gamma^* \rightarrow e^+ e^-$ .
- $X_1, X_2$ : two particles in thermalized medium.  $\gamma^*$ : virtual photon
- The rate per unit volume at finite temperature for this process is

$$\frac{d\Gamma}{d^4q} = \frac{e^2}{q^4} \int \frac{d^3l_+}{(2\pi)^3 2E_+} \frac{d^3l_-}{(2\pi)^3 2E_-} \delta^{(4)}(q - l_+ - l_-) W^{\mu\nu} L_{\mu\nu},$$

where  $E_{\pm} = \sqrt{|\mathbf{l}_{\pm}|^2 + m^2}$ ,  $e$  the electron charge and  $L_{\mu\nu}$  the lepton tensor

$$L^{\mu\nu} = 2(-q^2 g^{\mu\nu} + q^\mu q^\nu - \Delta l^\mu \Delta l^\nu).$$

- $q^\mu = l_+^\mu + l_-^\mu$ : virtual photon momentum.  $\Delta l^\mu = l_+^\mu - l_-^\mu$
- $L_{\mu\nu}$ : lepton tensor;  $W^{\mu\nu}$ : thermal averaged current correlator.

## Dilepton production rate contd...

$$\frac{d\Gamma}{d^4q} = \frac{e^2}{q^4} \int \frac{d^3l_+}{(2\pi)^3 2E_+} \frac{d^3l_-}{(2\pi)^3 2E_-} \delta^{(4)}(q - l_+ - l_-) W^{\mu\nu} L_{\mu\nu},$$

where  $E_{\pm} = \sqrt{|\mathbf{l}_{\pm}|^2 + m^2}$ ,  $e$  the electron charge and  $L_{\mu\nu}$  the lepton tensor

$$L^{\mu\nu} = 2(-q^2 g^{\mu\nu} + q^\mu q^\nu - \Delta l^\mu \Delta l^\nu).$$

$$W^{\mu\nu} = \langle w^{\mu\nu} \rangle,$$

where  $w^{\mu\nu}$  is given below and we have introduced the notation

$$\langle A \rangle = \int \frac{d^3p_1}{(2\pi)^3 2E_1} \frac{d^3p_2}{(2\pi)^3 2E_2} (2\pi)^4 \delta^{(4)}(q - p_1 - p_2) \frac{1}{e^{(u \cdot p_1)/T} \pm 1} \frac{1}{e^{(u \cdot p_2)/T} \pm 1} A.$$

## Dilepton production rate contd...

$$W^{\mu\nu} = W_1 \left( g^{\mu\nu} - \frac{q^\mu q^\nu}{q^2} \right) + W_2 \left( u^\mu - \frac{u \cdot q}{q^2} q^\mu \right) \left( u^\nu - \frac{u \cdot q}{q^2} q^\nu \right),$$

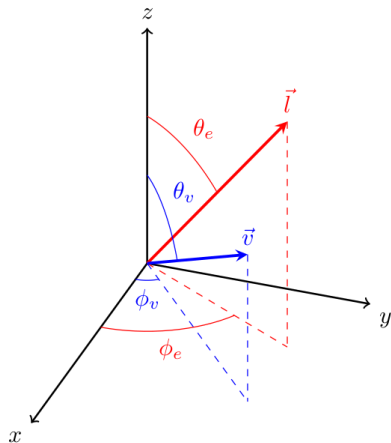
$$W_1 = \frac{\alpha a - \beta}{2a}, \quad W_2 = \frac{3\beta - \alpha a}{2a^2},$$

$$\alpha \equiv g_{\mu\nu} W^{\mu\nu}, \quad \beta \equiv u_\mu u_\nu W^{\mu\nu}, \quad a \equiv 1 - \frac{(u \cdot q)^2}{q^2}.$$

$$\frac{d\Gamma}{d^4q d\Omega_e} \propto \frac{1}{q^4} W_{\mu\nu} L^{\mu\nu}.$$

$$\frac{d\Gamma}{d^4q d\Omega_e} \propto \frac{1}{q^4} (-3q^2 - \Delta l^2) W_1 + \frac{1}{q^4} [-q^2 + (u \cdot q)^2 - (u \cdot \Delta l)^2] W_2.$$

# Relevant angles



$$u^\mu = \gamma(1, |\mathbf{v}| \sin \theta_v \cos \phi_v, |\mathbf{v}| \sin \theta_v \sin \phi_v, |\mathbf{v}| \cos \theta_v),$$
$$\Delta l^\mu = (0, 2|\mathbf{l}| \sin \theta_e \cos \phi_e, 2|\mathbf{l}| \sin \theta_e \sin \phi_e, 2|\mathbf{l}| \cos \theta_e),$$

# Anisotropic rate

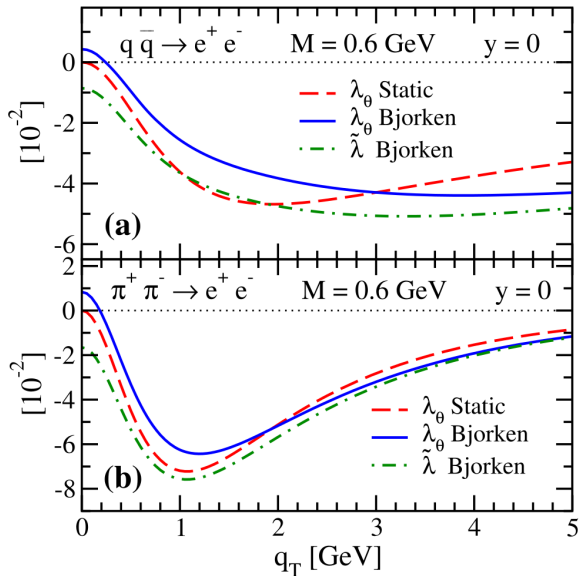
$$\frac{d\sigma}{d\Omega_e} \propto \mathcal{N}(1 + \lambda_\theta \cos^2 \theta_e + \lambda_\phi \sin^2 \theta_e \cos 2\phi_e + \lambda_{\theta\phi} \sin 2\theta_e \cos \phi_e + \lambda_\phi^\perp \sin^2 \theta_e \sin 2\phi_e + \lambda_{\theta\phi}^\perp \sin 2\theta_e \sin \phi_e),$$

$$\lambda_i = -W_2 4\gamma^2 |\mathbf{v}|^2 |\mathbf{l}|^2 \lambda'_i,$$

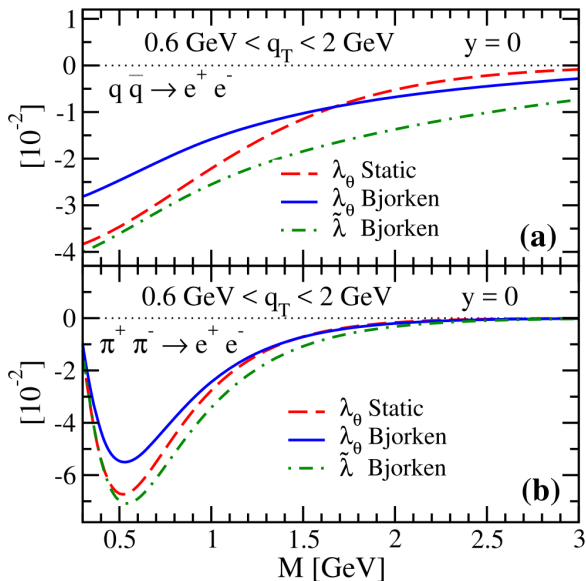
$$\mathcal{N} = (-3q^2 - \Delta l^2)W_1 + [-u^2 q^2 + (u \cdot q)^2 - 4\gamma^2 |\mathbf{v}|^2 |\mathbf{l}|^2 \lambda'_0]W_2,$$

$$\tilde{\lambda} \equiv \frac{\lambda_\theta + 3\lambda_\phi}{1 - \lambda_\phi}.$$

# Results



# Results contd...



- Virtual photon polarization and dilepton anisotropy in relativistic nucleus-nucleus collisions: [E. Speranza, A. Jaiswal, B. Friman, *Phys.Lett.B* 782 (2018) 395-400]. (Presented.)
- Virtual photon polarization in anisotropic medium (momentum anisotropy): [G. Baym, T. Hatsuda, M. Strickland, *Phys. Rev. C* 95 (2017) 044907].
- Polarization of direct photons from gluon anisotropy: [G. Baym, T. Hatsuda, *PTEP* 2015 (2015) 3, 031D01].



# Thank you!

